

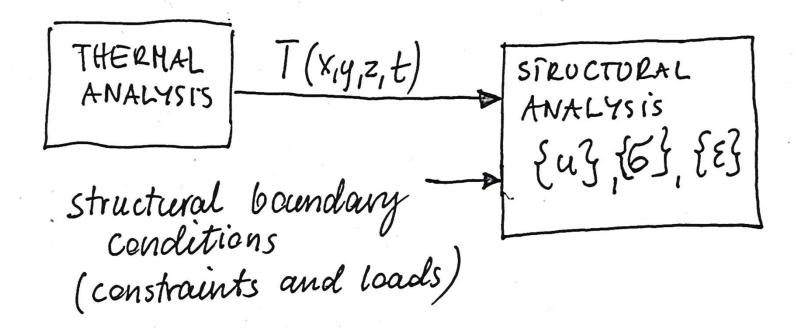
Institute of Aeronautics and Applied Mechanics

# Finite element method 2 (FEM 2)

Structural analysis with thermal effect

10.2021

## THERMAL STRESS. ANALYSIS



### STRUCTURAL ANALYSIS WITH THERMAL EFFECT

#### TOTAL STRAIN:

$$\{\xi\} = \{\xi\}_{\Gamma} + \{\xi\}_{e}$$
6×1
6×1

$$\{ \mathcal{E} \} = [B] \cdot \{ 9 \}_{e}$$

$$6 \times 1 \quad 6 \times n_{e} \quad n_{e} \times 1$$

ne-no. cf degrees of freedom in a finite element

STRESS

matrix

$$= [D] ([B] \cdot \{q\}_e - \{\xi\}_T)$$

$$= 6 \times 6 \cdot 6 \times Ne \text{ ne r1}$$

Virtual work principle:

Leggle 
$$\cdot$$
  $\{F\}_e - \sum_{1 \times 6} \{SE\} \cdot \{S\}_d \} dS_e = 0$ 

Then  $e \times 1$ 
 $\{SE\} = \{S\}_e \cdot \{S\}_e \}$ 

Then  $e \times 6$ 

Then  $e \times 6$ 

hence:

LSqJe 
$$\left(\begin{cases} F \rbrace_e - \int \begin{bmatrix} B \end{bmatrix}^T \cdot \left\{ 6 \right\} d\Omega_e \right) = 0$$

Ne ne x1

 $= 0$ 

$$\begin{cases} F_{e}^{2} = \int \begin{bmatrix} B \end{bmatrix}^{T} \cdot \{ 6 \} d\Omega e = \int \begin{bmatrix} B \end{bmatrix}^{T} \cdot [D] \left( \begin{bmatrix} B \end{bmatrix} \cdot \{ q \} - \{ E \} \} d\Omega e \\ n_{e} \times 1 & \text{Se} & n_{e} \times 6 & 6 \times 6 \end{cases}$$

$$\begin{cases} F_{e}^{2} = \int \begin{bmatrix} B \end{bmatrix}^{T} \cdot \{ 6 \} d\Omega e = \int \begin{bmatrix} B \end{bmatrix}^{T} \cdot [D] \left( \begin{bmatrix} B \end{bmatrix} \cdot \{ q \} - \{ E \} \} d\Omega e \\ n_{e} \times 1 & 6 \times 1 \end{cases}$$

$$\begin{cases} F_{e}^{2} = \int \begin{bmatrix} B \end{bmatrix}^{T} \cdot \{ 6 \} d\Omega e = \int \begin{bmatrix} B \end{bmatrix}^{T} \cdot [D] \left( \begin{bmatrix} B \end{bmatrix} \cdot \{ q \} - \{ E \} \} d\Omega e \\ n_{e} \times 1 & 6 \times 1 \end{cases}$$

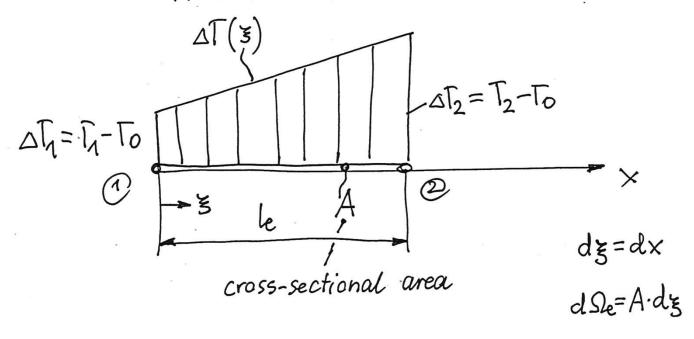
ne x ne

local stiffness matrix

vector of thermal load

Finally:  $[k]_{e} \cdot \{q\}_{e} = \{F\}_{e} + \{F_{T}\}_{e}$   $n_{e \times n_{e}} = n_{e \times 1} = n_{e \times 1}$ hodal forces due to temperature Set of equations for the entire FE model:  $[K], \{q\} = \{F\} + \{F_{+}\}$ NOOF \* NOOF \* NOOF \* L NDOF×1 global vector of structural of nodal para" global vector of thermal load global load vector parameters NDOF - number of degrees of freedom

EXAMPLE: FIND THERMAL LOAD VECTOR FOR A BAR FINITE ELEMENT.



ΔT, ΔT2 - nodal parameters in thermal analysis

$$\Delta T(\xi) = \lfloor N_1(\xi), N_2(\xi) \rfloor \cdot \begin{cases} \Delta I_1 \\ \Delta I_2 \end{cases}$$

$$N_1(\xi) = 1 - \frac{\xi}{le}$$
 $N_2(\xi) = \frac{\xi}{le}$ 

$$\mathcal{E}_{T}(\mathfrak{F}) = \infty_{se} \cdot \Delta \overline{\Gamma}(\mathfrak{F}) = \infty_{se} \cdot [N_{1}, N_{2}] \cdot \{\Delta \overline{\Gamma}_{2}\}$$

$$for \\ "3D": [B] = [R] \cdot [N]$$

$$6 \times ne \quad 6 \times 3 \quad 3 \times ne$$

$$for \\ 1D": [B] = \frac{\partial}{\partial \mathfrak{F}} [N_{1}, N_{2}] = [\frac{\partial N_{1}}{\partial \mathfrak{F}}, \frac{\partial N_{2}}{\partial \mathfrak{F}}] = [-\frac{1}{le}, \frac{1}{le}]$$

$$\begin{cases} B \\ = \\ \begin{cases} -\frac{1}{le} \\ \\ le \end{cases} \end{cases}$$

(FOR 
$$3D''$$
  $\{6\} = [D] \cdot \{\xi\}_e$ ) here:  $6 = E \cdot \xi_e$  (along the bar axis)

THERMAL LOAD VECTOR:

$$\begin{cases}
F_{\tau} \\ e = \int \{B\} \cdot E \cdot \mathcal{E}_{\tau}(\xi) d\Omega_{e} = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{1}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot \{\Delta \overline{L}_{2}\} \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N_{2}(\xi)] \cdot A \cdot d\xi = \\
= \int \{-\frac{1}{le}\} E \cdot \alpha_{se} \cdot [N_{\eta}(\xi), N$$

$$= \begin{cases} F_{\tau} = \alpha_{se} EA \\ 2 \times 4 \end{cases} = \begin{cases} \frac{(1 - \frac{\xi}{le})}{le} - \frac{\xi}{le^2} \\ \frac{(1 - \frac{\xi}{le})}{le} - \frac{\xi}{le^2} \\ \frac{(1 - \frac{\xi}{le})}{le} - \frac{\xi}{le^2} \end{cases} = \alpha_{se} EA \begin{cases} -\frac{1}{2} - \frac{1}{2} \\ \frac{1}{2} - \frac{1}{2} \end{bmatrix} \cdot \begin{cases} \Delta \Gamma_1 \\ \Delta \Gamma_2 \end{cases} = \frac{\Delta \Gamma_1 + \Delta \Gamma_2}{2} \quad \alpha_{se} EA \cdot \begin{cases} -1 \\ 1 \end{cases}$$